



On a maximal subgroup of the orthogonal group $O_8^+(3)$

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Abstract

The orthogonal simple group $O_8^+(3)$ has three conjugacy classes of maximal subgroups of the form $3^6:L_4(3)$. These groups are all isomorphic to each other and each group has order 4421589120 with index 1120 in $O_8^+(3)$. In this paper, we will compute the ordinary character table of one of these classes of maximal subgroups using the technique of Fischer-Clifford matrices. This technique is very efficient to compute the ordinary character table of an extension group $\overline{G} = N.G$ and especially where the normal subgroup N of \overline{G} is an elementary abelian p -group. The said technique reduces the computation of the ordinary character table of \overline{G} to find a handful of so-called Fischer-Clifford matrices of \overline{G} and the ordinary or projective character tables of the inertia factor groups of the action of \overline{G} on N .

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1. Introduction

The orthogonal group $O_8^+(3)$ of order $4952179814400 = 2^{12} \cdot 3^{12} \cdot 5^2 \cdot 7 \cdot 13$ has 27 conjugacy classes of maximal subgroups [8]. Among the maximal subgroups of $O_8^+(3)$ are three non-conjugate but isomorphic subgroups \overline{G}_1 , \overline{G}_2 and \overline{G}_3 of the form $3^6:L_4(3)$ with order 4421589120 and index 1120 in $O_8^+(3)$. The aim of this paper is to compute the Fischer-Clifford matrices and hence the character table of \overline{G}_1 , the first group of the three as they appear in the ATLAS [8]. For this purpose, the Fischer-Clifford matrices technique which is based on Clifford theory and was developed by Bernd Fischer [9] is used. The group $3^6:L_4(3)$ which we shall now denote by \overline{G} is a split-extension of $N = 3^6$, the vector space of dimension 6 over $GF(3)$, by the linear group $G = L_4(3) \cong O_6^+(3)$.

Let $\overline{G} = N:G$ be a split extension of N by G , where N is a vector space of dimension n over $GF(p)$, for a prime p , on which a linear group G acts naturally. The Fischer-Clifford technique involves the construction of a non-singular matrix $M(g)$ for each conjugacy class representative g of G , which together with the fusion maps and ordinary character tables of some subgroups of G , called the inertia factor groups, are used to assemble the complete ordinary character table of \overline{G} .

The Fischer-Clifford matrix $M(g)$ is partitioned row-wise into blocks, where each block corresponds to an inertia group \overline{H}_i of $\theta_i \in \text{Irr}(N)$ in \overline{G} . Using the columns of the character tables of the inertia factors $H_i \cong \frac{\overline{H}_i}{N}$ which correspond to classes of H_i which fuse to the class $[g]$ of G and multiplying these columns by the rows of the Fischer-Clifford matrix $M(g)$ that correspond to H_i , a portion of the character table of \overline{G} which is in the block corresponding to \overline{H}_i for the classes that come from the coset Ng is constructed. The character table of \overline{G} is thus divided row-wise into blocks, where each block corresponds to an inertia group $\overline{H}_i = N:H_i$. The reader is referred to [2], [15], [16], [19], [20] and [21] for more literature on this technique. A brief theoretical background of the Fischer-Clifford theory is given in Section 2.

In Section 3, the coset analysis technique [17] is used to determine the conjugacy classes of \overline{G} . In Sections 4 and 5 the inertia factor groups H_i and their fusion maps into $G = L_3(4)$ are computed. The Fischer-Clifford matrices of $\overline{G} = 3^6:L_4(3)$ are determined in Section 6 and the associated

ordinary character table of \overline{G} is to be found in Section 7. The technique of set intersection of characters (see [1], [17], [18]) is mainly used to compute the fusion of the conjugacy classes of $\overline{G} = 3^6:L_4(3)$ into $O_8^+(3)$. Most of our computations are carried out with the computer algebra systems MAGMA [7] and GAP [11] and the notation of ATLAS is mostly followed.

2. Theory of Fischer-Clifford Matrices

Let $\overline{G} = N:G$ be a split extension of N by G . Then for $\theta \in \text{Irr}(N)$, we define $\overline{H} = \{x \in \overline{G} | \theta^x = \theta\} = I_{\overline{G}}(\theta)$ and $H = \{x \in G | \theta^x = \theta\} = I_G(\theta)$ where $I_{\overline{G}}(\theta)$ is the stabilizer of θ in the action of \overline{G} on $\text{Irr}(N)$, we have that $I_{\overline{G}}(\theta)$ is a subgroup of \overline{G} and N is normal subgroup in $I_{\overline{G}}(\theta)$. Also $[\overline{G}:I_{\overline{G}}(\theta)]$ is the size of the orbit containing θ . Then it can be shown that $\overline{H} = N:H$, where \overline{H} is the inertia group of θ in \overline{G} . The inertia factor $\overline{H}/N \cong H$ can be regarded as the inertia group of θ in the factor group $\overline{G}/N \cong G$. Define θ^g by $\theta^g(n) = \theta(gng^{-1})$ for $g \in \overline{G}$, $n \in N$, then $\theta^g \in \text{Irr}(N)$. We say that θ is extendible to \overline{H} if there exists $\varphi \in \text{Irr}(\overline{H})$ such that $\varphi \downarrow N = \theta$. If θ is extendible to \overline{H} then by Gallagher [10], we have $\{\varphi | \varphi \in \text{Irr}(\overline{H}), \varphi \downarrow N = \theta\} = \{\overline{\beta}\varphi | \beta \in \text{Irr}(\overline{H}/N)\}$, where $\overline{\beta} \in \text{Irr}(\overline{H})$ is a lifting for β into \overline{H} . Let \overline{G} have the property that every irreducible character of N can be extended to its inertia group. Now let $\theta_1 = 1_N, \theta_2, \dots, \theta_t$ be representatives of the orbits of \overline{G} on $\text{Irr}(N)$, $\overline{H}_i = I_{\overline{G}}(\theta_i)$, $1 \leq i \leq t$, $\varphi_i \in \text{Irr}(\overline{H}_i)$ be an extension of θ_i to \overline{H}_i and $\overline{\beta} \in \text{Irr}(\overline{H}_i)$ such that $N \subseteq \text{Ker}(\overline{\beta})$. Then it can be shown that

$$\begin{aligned} \text{Irr}(\overline{G}) &= \bigcup_{i=1}^t \{(\overline{\beta}\varphi_i)^{\overline{G}} | \overline{\beta} \in \text{Irr}(\overline{H}_i), N \subseteq \text{Ker}(\overline{\beta})\} \\ &= \bigcup_{i=1}^t \{(\overline{\beta}\varphi_i)^{\overline{G}} | \beta \in \text{Irr}(\overline{H}_i/N)\} \end{aligned}$$

Hence the irreducible characters of \overline{G} will be divided into blocks, where each block corresponds to an inertia group \overline{H}_i . Let H_i be the inertia factor group and φ_i be an extension of θ_i to \overline{H}_i . Take $\theta_1 = 1_N$ as the identity character of N , then $\overline{H}_1 = \overline{G}$ and $H_1 \cong G$. Let $X(g) = \{x_1, x_2, \dots, x_{c(g)}\}$ be a set of representatives of the conjugacy classes of \overline{G} from the coset $N\overline{g}$ whose images under the natural homomorphism $\overline{G} \rightarrow G$ are in $[g]$ and we take $x_1 = \overline{g}$. We define,

$$R(g) = \{(i, y_k) | 1 \leq i \leq t, H_i \cap [g] \neq \emptyset, 1 \leq k \leq r\}$$

and we note that y_k runs over representatives of the conjugacy classes of elements of H_i which fuse into $[g]$ in G . Then we define the Fischer-Clifford

matrix $M(g)$ by $M(g) = (a_{(i,y_k)}^j)$, where $a_{(i,y_k)}^j = \sum_l' \frac{|C_{\overline{G}}(x_j)|}{|C_{\overline{H}_i}(y_{l_k})|} \varphi_i(y_{l_k})$ with columns indexed by $X(g)$ and rows indexed by $R(g)$ and where \sum_l' is the summation over all l for which $y_{l_k} \sim x_j$ in \overline{G} . Then the partial character table of \overline{G} on the classes $\{x_1, x_2, \dots, x_{c(g)}\}$ is given by

$$\begin{pmatrix} C_1(g)M_1(g) \\ C_2(g)M_2(g) \\ \vdots \\ C_t(g)M_t(g) \end{pmatrix}$$

where the Fischer-Clifford matrix

$$M(g) = \begin{pmatrix} M_1(g) \\ M_2(g) \\ \vdots \\ M_t(g) \end{pmatrix}$$

is divided into blocks $M_i(g)$ with each block corresponding to an inertia group \overline{H}_i and $C_i(g)$ is the partial character table of H_i consisting of the columns corresponding to the classes that fuse into $[g]$ in G . We can also observe that the number of irreducible characters of \overline{G} is the sum of the number of irreducible characters of the inertia factors H_i 's. For complete information on the properties of Fischer-Clifford matrices the reader is referred to [2], [15], [16], [19], [20] and [21]. The group $\overline{G} = 3^6:L_4(3)$ is a split extension with 3^6 abelian and therefore by Mackey's theorem (see Theorem 5.1.15 in [18]), we have that each irreducible character of 3^6 can be extended to its inertia group in \overline{G} . With this theoretical assertion in mind, the character table of $\overline{G} = 3^6:L_4(3)$ is now going to be determined using the above outline.

3. The Conjugacy Classes of $\overline{G} = 3^6:L_4(3)$

In this section, the method of coset analysis is used to determine the conjugacy classes of the elements of $\overline{G} = 3^6:L_4(3)$. This method was developed and first used by Moori in [17] and since then, it has been used by many other researchers to compute the conjugacy classes of groups of extension type. The reader is referred to [3] and [4] for recent application of this technique. By making use of the standard generators from the online ATLAS of Group Representations [25], the groups $P = O_8^+(3)$

and \overline{G} are represented as permutations on 1080 points in MAGMA. The command “*IsMaximal*(P, \overline{G});” confirms that \overline{G} is a maximal subgroup of P . Proceeding with the commands, “ $a, b := \text{ChiefSeries}(\overline{G});$ ”, “ $N := a[2];$ ”, “ $M := \text{GModule}(\overline{G}, N);$ ” and “ $M:\text{Maximal};$ ” the group $G = L_4(3)$ is constructed as matrix group of degree 6 over $GF(3)$. Note that N is the only elementary abelian 3-group of order 729 in \overline{G} . The following two 6×6 matrices g_1 and g_2 of orders 2 and 12, respectively, are obtained as the generators of G .

$$g_1 = \begin{pmatrix} 0 & 1 & 0 & 0 & 0 & 0 \\ 1 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 1 & 0 \\ 0 & 0 & 2 & 1 & 1 & 0 \\ 0 & 0 & 1 & 0 & 0 & 0 \\ 1 & 2 & 1 & 0 & 2 & 1 \end{pmatrix} \quad g_2 = \begin{pmatrix} 0 & 0 & 1 & 0 & 0 & 0 \\ 0 & 0 & 0 & 1 & 0 & 0 \\ 2 & 1 & 1 & 1 & 0 & 0 \\ 2 & 0 & 0 & 1 & 1 & 0 \\ 0 & 0 & 0 & 0 & 0 & 1 \\ 1 & 0 & 1 & 0 & 1 & 0 \end{pmatrix}$$

The 29 conjugacy classes of $G = \langle g_1, g_2 \rangle$ are computed within GAP.

3.1. The Action of $G = L_4(3)$ on 3^6

Let $\overline{G} = 3^6:L_4(3)$ be the split extension of $G = \langle g_1, g_2 \rangle$ by $N = 3^6$, where N is considered as a vector space $V_6(3)$ of dimension 6 over $GF(3)$. Also, $N \cong V_6(3)$ as a G -module of $G = \langle g_1, g_2 \rangle$ is irreducible. Using GAP, it turns out that the action of $G = L_4(3)$ on $N = 3^6$ has four orbits of lengths 1, 234, 234 and 260 with corresponding point stabilizers P_1, P_2, P_3 and P_4 .

3.2. Permutation Character of $G = L_4(3)$ on 3^6

Checking the indices of maximal subgroups of $G = L_4(3)$ in the ATLAS [8], P_2 and P_3 sit maximally inside the maximal subgroups with the structure $U_4(2):2$ while P_4 sits maximal inside the maximal subgroup $3^4:2(A_4 \times A_4).2$ of G . It follows that $P_1 = L_4(3), P_2 = U_4(2), P_3 = U_4(2)$ and $P_4 = 3^4:2(A_4 \times A_4)$ of indices 1, 234, 234 and 260 respectively in $L_4(3)$. We will now determine, with the use of the permutation character $\chi(L_4(3)|3^6)$ of G on N whether P_2 and P_3 are sitting separately inside one of the two classes of $U_4(2):2$ or both are in one of the classes of $U_4(2):2$. The permutation character $\chi(L_4(3)|3^6) = 1 + I_{P_2}^{L_4(3)} + I_{P_3}^{L_4(3)} + I_{P_4}^{L_4(3)}$ of G acting on N is now going to be computed, where $I_{P_2}^{L_4(3)}, I_{P_3}^{L_4(3)}$ and $I_{P_4}^{L_4(3)}$ are the identity characters of P_2, P_3 and P_4 induced to G respectively. To determine

$I_{P_2}^{L_4(3)}$, the fusion of conjugacy classes of P_2 into $L_4(3)$ and the restrictions of $\chi_i \in \text{Irr}(L_4(3))$ to P_2 , where $\deg(\chi_i) < 234$ are used. Thus restricting $\chi_i \in \text{Irr}(L_4(3))$ to P_2 where $i \in 1, 2, 3, \dots, 8$, and computing the inner product $\langle \chi_i, \psi_1 \rangle$ of each χ_i , $i \in 1, 2, 3, \dots, 8$, with the identity character ψ_1 of P_2 , the values below are obtained.

	χ_1	χ_2	χ_3	χ_4	χ_5	χ_6	χ_7	χ_8
$\langle \chi_i, \psi_1 \rangle$	1	0	1	0	1	0	1	1

From the above table and taking into consideration the Frobenius-Reciprocity theorem [12], the permutation character $I_{P_2}^{L_4(3)} = 1a + 26b + 52a + 65b + 90a$ is obtained. Similarly, $I_{P_4}^{L_4(3)}$ is determined. In this case, we restrict $\chi_i \in \text{Irr}(L_4(3))$, where $i \in 1, 2, 3, \dots, 10$ to P_4 and let ψ_1 be the identity character P_4 . Computing the inner product $\langle \chi_i, \psi_1 \rangle$ of each χ_i for $i \in 1, 2, 3, \dots, 10$ with ψ_1 , the values below are found.

	χ_1	χ_2	χ_3	χ_4	χ_5	χ_6	χ_7	χ_8	χ_9	χ_{10}
$\langle \chi_i, \psi_1 \rangle$	1	0	0	1	0	1	1	1	0	0

The identity character $I_{P_4}^{L_4(3)}$ of P_4 induced to $L_4(3)$ is therefore given by, $I_{P_4}^{L_4(3)} = 1a + 39a + 65a + 65b + 90a$. It follows that the permutation character $\chi(L_4(3)|3^6)$ is given as,

$$\begin{aligned}
 \chi(L_4(3)|3^6) &= 1 + 2I_{P_2}^{L_4(3)} + I_{P_4}^{L_4(3)} \\
 &= 1a + 2(1a + 26b + 52a + 65b + 90a) + 1a + 39a + 65a + 65b + 90a \\
 &= 4 \times 1a + 2 \times 26b + 39a + 2 \times 52a + 65a + 3 \times 65b + 3 \times 90a.
 \end{aligned}$$

The permutation characters $\chi(L_4(3)|P_i)$ are written in terms of the ordinary irreducible characters of G and are computed directly using the character table of G . The permutation character $\chi(L_4(3)|3^6)$ on the different conjugacy classes of G determines the number k of fixed points of each $g \in G$ in 3^6 . The values of k obtained by the above permutation character are listed in Table 1.

Table 1: Permutation Character of $G = L_4(3)$ on 3^6

$[g]_G$	1A	2A	2B	3A	3B	3C	3D	4A	4B	4C	5A	6A	6B	6C	6D
$\chi(G P_1)$	1	1	1	1	1	1	1	1	1	1	1	1	1	1	1
$\chi(G P_2)$	234	30	2	18	0	18	36	0	2	4	4	0	6	2	2
$\chi(G P_3)$	234	30	2	18	0	18	36	0	2	4	4	0	6	2	2
$\chi(G P_4)$	260	20	4	44	8	26	26	0	4	0	0	2	2	4	4
k	729	81	9	81	9	63	99	1	9	9	9	3	15	9	9
	6E	8A	9A	9B	10A	12A	12B	12C	13A	13B	13C	13D	20A	20B	
$\chi(G P_1)$	1	1	1	1	1	1	1	1	1	1	1	1	1	1	
$\chi(G P_2)$	2	0	0	6	0	0	0	2	0	0	0	0	0	0	
$\chi(G P_3)$	2	0	0	6	0	0	0	2	0	0	0	0	0	0	
$\chi(G P_4)$	4	0	2	2	0	0	0	4	0	0	0	0	0	0	
k	9	1	3	15	1	1	1	9	1	1	1	1	1	1	

Clearly, $\chi(L_4(3)|3^6) = 4 \times 1a + 2 \times 26b + 39a + 2 \times 52a + 65a + 3 \times 65b + 3 \times 90a$ is not the required permutation character of G acting on N since $k \neq 3^n$, $n \in \{0, 1, 2, 3, 4, 5, 6\}$ for all the conjugacy classes of G . In order to obtain values of k such that $k = 3^n$, $n \in \{0, 1, 2, 3, 4, 5, 6\}$, another possible fusion of conjugacy classes of $P_2 \cong P_3$ into $L_4(3)$ is considered. We restrict $\chi_i \in Irr(L_4(3))$ to P_2 , and then compute the inner product $\langle \chi_i, \psi_1 \rangle$ of each χ_i , $i \in 1, 2, 3, \dots, 8$ with the identity character ψ_1 of P_2 . The values of $\langle \chi_i, \psi_1 \rangle$ are listed below,

	χ_1	χ_2	χ_3	χ_4	χ_5	χ_6	χ_7	χ_8
$\langle \chi_i, \psi_1 \rangle$	1	1	0	0	1	1	0	1

From the above table and taking into consideration the Frobenius-Reciprocity theorem, the permutation character $I_{P_2}^{L_4(3)}$ now assumes the following form,

$$I_{P_2}^{L_4(3)} = 1a + 26a + 52a + 65a + 90a.$$

Using $I_{P_2}^{L_4(3)} = 1a + 26a + 52a + 65a + 90a$ and $I_{P_3}^{L_4(3)} = 1a + 26b + 52a + 65b + 90a$ we obtain that

$$\begin{aligned}
\chi(L_4(3)|3^6) &= 1 + I_{P_2}^{L_4(3)} + I_{P_3}^{L_4(3)} + I_{P_4}^{L_4(3)} \\
&= 4 \times 1a + 1 \times 26a + 1 \times 26b + 1 \times 39a + 2 \times 52a \\
&\quad + 2 \times 65a + 2 \times 65b + 3 \times 90a.
\end{aligned}$$

Using this result, the correct number k of fixed points of each $g \in G$ in 3^6 is obtained and is listed in Table 2. This confirms that P_2 and P_3 sit separately inside the two maximal subgroups of G of the form $U_4(2):2$.

Table 2: Permutation Character of $G = L_4(3)$ on 3^6

$g _G$	1A	2A	2B	3A	3B	3C	3D	4A	4B	4C	5A	6A	6B	6C	6D
$\chi(G P_1)$	1	1	1	1	1	1	1	1	1	1	1	1	1	1	1
$\chi(G P_2)$	234	30	2	18	0	18	36	0	2	4	4	0	6	2	2
$\chi(G P_3)$	234	30	2	18	0	36	18	0	2	4	4	6	0	2	2
$\chi(G P_4)$	260	20	4	44	8	26	26	0	4	0	0	2	2	4	4
k	729	81	9	81	9	81	81	1	9	9	9	9	9	9	9
$g _G$	6E	8A	9A	9B	10A	12A	12B	12C	13A	13B	13C	13D	20A	20B	
$\chi(G P_1)$	1	1	1	1	1	1	1	1	1	1	1	1	1	1	
$\chi(G P_2)$	2	0	0	6	0	0	0	2	0	0	0	0	0	0	
$\chi(G P_3)$	2	0	6	0	0	0	0	2	0	0	0	0	0	0	
$\chi(G P_4)$	4	0	2	2	0	0	0	4	0	0	0	0	0	0	
k	9	1	9	9	1	1	1	9	1	1	1	1	1	1	

3.3. The Classes of $\overline{G} = 3^6:L_4(3)$

The values of k enable us to determine the number f_j of orbits Q_i 's, $1 \leq i \leq k$ that fuse together under the action of $C_G(g)$ to form one orbit Δ_j (see [18]). To determine the values of these f_j 's and the orders of class representatives $dg \in \overline{G}$, Programmes A and B in [24] written in GAP are used, respectively. If $o(g) = m$ and $w = 1_N$ then $o(dg) = m$ and if $w \neq 1_N$ then $o(dg) = 3m$ (see Theorem 2.3.10 in [18]). The formula $|C_{\overline{G}}(x)| = \frac{k}{f_j}|C_G(g)|$ is then used to calculate the order of the centralizer of each class of \overline{G} with representative x and a constant $m_j = \frac{f_j}{k}|N|$ is also calculated for each value of f_j . This constant plays a very crucial role in determination of the entries of the Fischer Clifford matrices. The group $\overline{G} = 3^6:L_4(3)$ is found to have 111 conjugacy classes of elements. Table 3 below gives a detailed information on the conjugacy classes of $\overline{G} = 3^6:L_4(3)$. The power maps of elements of \overline{G} are given in the second last column of Table 3 whereas the fusion of \overline{G} into $O_8^+(3)$, as determined in Section 8 of this paper, is found in the last column of Table 3.

Table 3: The Conguajacy Classes of $\overline{G} = 3^6:L_4(3)$

$[g]_G$	k	f_j	m_j	d_j	w	$[x]_{\overline{G}}$	$ C_{\overline{G}}(x) $	2	3	5	7	13	$\mapsto O_8^+(3)$
1A	729	1	1	(0, 0, 0, 0, 0, 0)	(0, 0, 0, 0, 0, 0)	1A	4421589120						1A
		234	234	(0, 0, 0, 1, 0, 1)	(0, 0, 0, 1, 0, 1)	3A	18895680		1A				3B
		234	234	(0, 0, 0, 1, 0, 2)	(0, 0, 0, 1, 0, 2)	3B	18895680		1A				3E
		260	260	(0, 0, 0, 0, 0, 1)	(0, 0, 0, 0, 0, 1)	3C	17006112		1A				3A
2A	81	1	9	(0, 0, 0, 0, 0, 0)	(0, 0, 0, 0, 0, 0)	2A	233280	1A					2A
		20	180	(0, 0, 0, 0, 0, 1)	(0, 0, 0, 0, 0, 2)	6A	11664	3C	2A				6A
		30	270	(0, 0, 0, 0, 1, 0)	(1, 1, 1, 2, 1, 2)	6B	7776	3A	2A				6E
		30	270	(0, 0, 0, 1, 0, 2)	(0, 0, 0, 2, 0, 1)	6C	7776	3B	2A				6H
2B	9	1	81	(0, 0, 0, 0, 0, 0)	(0, 0, 0, 0, 0, 0)	2B	10368	1A					2D
		2	162	(0, 0, 0, 0, 0, 1)	(2, 1, 2, 1, 2, 1)	6D	5184	3B	2B				6N
		2	162	(0, 0, 0, 0, 1, 1)	(0, 0, 2, 1, 2, 2)	6E	5184	3A	2B				6K
		4	324	(0, 0, 0, 0, 1, 0)	(1, 2, 0, 0, 0, 1)	6F	2592	3C	2B				6D
3A	81	1	9	(0, 0, 0, 0, 0, 0)	(0, 0, 0, 0, 0, 0)	3D	472392		1A				3A
		1	9	(0, 0, 0, 1, 0, 1)	(0, 0, 0, 0, 0, 0)	3E	472392		1A				3K
		1	9	(0, 0, 0, 2, 0, 2)	(0, 0, 0, 0, 0, 0)	3F	472392		1A				3J
		1	9	(0, 1, 0, 0, 2, 0)	(0, 0, 0, 0, 0, 0)	3G	472392		1A				3H
		1	9	(0, 1, 0, 1, 2, 1)	(0, 0, 0, 0, 0, 0)	3H	472392		1A				3G
		1	9	(0, 1, 0, 2, 2, 2)	(0, 0, 0, 0, 0, 0)	3I	472392		1A				3F
		1	9	(0, 2, 0, 0, 1, 0)	(0, 0, 0, 0, 0, 0)	3J	472392		1A				3I
		1	9	(0, 1, 0, 1, 1, 1)	(0, 0, 0, 0, 0, 0)	3K	472392		1A				3C
		1	9	(0, 2, 0, 2, 1, 2)	(0, 0, 0, 0, 0, 0)	3L	472392		1A				3D
		72	648	(0, 0, 0, 0, 0, 1)	(0, 0, 0, 0, 0, 0)	3M	6561		1A				3M
3B	9	1	81	(0, 0, 0, 0, 0, 0)	(0, 0, 0, 0, 0, 0)	3N	729		1A				3M
		1	81	(0, 0, 0, 0, 0, 1)	(2, 2, 0, 1, 0, 1)	9A	729		3C				9E
		1	81	(0, 0, 0, 0, 0, 2)	(1, 1, 0, 2, 0, 2)	9B	729		3C				9B
		1	81	(0, 0, 0, 0, 1, 0)	(0, 1, 0, 2, 2, 1)	9C	729		3C				9H
		1	81	(0, 0, 0, 0, 1, 1)	(2, 0, 0, 0, 2, 2)	9D	729		3C				9C
		1	81	(0, 0, 0, 0, 1, 2)	(1, 2, 0, 1, 2, 0)	9E	729		3C				9I
		1	81	(0, 0, 0, 0, 2, 0)	(0, 2, 0, 1, 1, 2)	9F	729		3C				9G
		1	81	(0, 0, 0, 0, 2, 1)	(2, 1, 0, 2, 1, 0)	9G	729		3C				9J
		1	81	(0, 0, 0, 0, 2, 2)	(1, 0, 0, 0, 1, 1)	9H	729		3C				9F
3C	81	1	9	(0, 0, 0, 0, 0, 0)	(0, 0, 0, 0, 0, 0)	3O	157464		1A				3B
		4	36	(0, 0, 0, 0, 1, 0)	(0, 0, 0, 0, 0, 0)	3P	39366		1A				3H
		4	36	(0, 0, 0, 0, 2, 0)	(0, 0, 0, 0, 0, 0)	3Q	39366		1A				3I
		6	54	(0, 0, 0, 1, 2, 1)	(0, 0, 0, 0, 0, 0)	3R	26244		1A				3L
		12	108	(0, 0, 0, 1, 1, 1)	(0, 0, 0, 0, 0, 0)	3S	13122		1A				3M
		54	486	(0, 0, 0, 0, 0, 1)	(2, 2, 0, 1, 0, 1)	9I	2916		3C				9A
3D	81	1	9	(0, 0, 0, 0, 0, 0)	(0, 0, 0, 0, 0, 0)	3T	157464		1A				3E
		4	36	(0, 0, 1, 0, 0, 0)	(0, 0, 0, 0, 0, 0)	3U	39366		1A				3J
		4	36	(0, 0, 1, 1, 2, 1)	(0, 0, 0, 0, 0, 0)	3V	39366		1A				3K
		6	54	(0, 0, 0, 1, 0, 1)	(0, 0, 0, 0, 0, 0)	3W	26244		1A				3L
		12	108	(0, 0, 1, 2, 1, 2)	(0, 0, 0, 0, 0, 0)	3X	13122		1A				3M
		54	486	(0, 0, 0, 0, 0, 1)	(1, 1, 0, 2, 0, 2)	9J	2916		3C				9D
4A	1	1	729	(0, 0, 0, 0, 0, 0)	(0, 0, 0, 0, 0, 0)	4A	1440	2A					4B
4B	9	1	81	(0, 0, 0, 0, 0, 0)	(0, 0, 0, 0, 0, 0)	4B	864	2B					4A
		2	162	(0, 0, 0, 0, 0, 1)	(1, 2, 1, 2, 1, 2)	12A	432	6D	4B				12G
		2	162	(0, 0, 0, 0, 1, 0)	(2, 1, 0, 0, 0, 2)	12B	432	6F	4B				12B
		2	162	(0, 0, 0, 0, 1, 1)	(0, 0, 1, 2, 1, 1)	12C	432	6E	4B				12D
		2	162	(0, 0, 0, 0, 1, 2)	(1, 2, 2, 1, 2, 0)	12D	432	6F	4B				12C
4C	9	1	81	(0, 0, 0, 0, 0, 0)	(0, 0, 0, 0, 0, 0)	4C	288	2A					4B
		4	324	(0, 0, 0, 0, 0, 1)	(1, 0, 1, 1, 2, 0)	12E	72	6C	4C				12R
		4	324	(0, 0, 0, 0, 1, 0)	(0, 1, 2, 1, 2, 1)	12F	72	6B	4C				12O
5A	9	1	81	(0, 0, 0, 0, 0, 0)	(0, 0, 0, 0, 0, 0)	5A	180			1A			5A
		4	324	(0, 0, 0, 0, 1, 0)	(2, 2, 2, 1, 1, 1)	15A	45		5A	3A			15A
		4	324	(0, 0, 1, 0, 0, 0)	(1, 0, 2, 2, 2, 2)	15B	45		5A	3B			15D
6A	9	1	81	(0, 0, 0, 0, 0, 0)	(0, 0, 0, 0, 0, 0)	6G	648	3T	2A				6H
		2	162	(0, 0, 0, 1, 2, 1)	(0, 0, 0, 0, 0, 0)	6H	324	3W	2A				6Q
		6	486	(0, 0, 0, 1, 0, 2)	(2, 2, 0, 1, 0, 1)	18A	108	9J	6A				18D

Table 3: (continued)

$[g]_G$	k	f_j	m_j	d_j	w	$[x]_{\overline{G}}$	$ C_{\overline{G}}(x) $	2	3	5	7	13	$\mapsto O_8^+(3)$
$6B$	9	1	81	(0, 0, 0, 0, 0, 0)	(0, 0, 0, 0, 0, 0)	$6I$	648	$3O$	$2A$				$6E$
		2	162	(0, 0, 0, 1, 1, 1)	(0, 0, 0, 0, 0, 0)	$6J$	324	$3R$	$2A$				$6Q$
		6	486	(0, 1, 0, 0, 1, 0)	(1, 1, 0, 2, 0, 2)	$18B$	108	$9I$	$6A$				$18A$
$6C$	9	1	81	(0, 0, 0, 0, 0, 0)	(0, 0, 0, 0, 0, 0)	$6K$	648	$3D$	$2B$				$6D$
		1	81	(0, 0, 0, 0, 0, 1)	(0, 0, 0, 0, 0, 0)	$6L$	648	$3F$	$2B$				$6V$
		1	81	(0, 0, 0, 0, 0, 2)	(0, 0, 0, 0, 0, 0)	$6M$	648	$3E$	$2B$				$6W$
		1	81	(0, 0, 0, 0, 1, 0)	(0, 0, 0, 0, 0, 0)	$6N$	648	$3K$	$2B$				$6L$
		1	81	(0, 0, 0, 0, 1, 1)	(0, 0, 0, 0, 0, 0)	$6O$	648	$3J$	$2B$				$6U$
		1	81	(0, 0, 0, 0, 1, 2)	(0, 0, 0, 0, 0, 0)	$6P$	648	$3L$	$2B$				$6M$
		1	81	(0, 0, 0, 0, 2, 0)	(0, 0, 0, 0, 0, 0)	$6Q$	648	$3I$	$2B$				$6O$
		1	81	(0, 0, 0, 0, 2, 1)	(0, 0, 0, 0, 0, 0)	$6R$	648	$3H$	$2B$				$6P$
		1	81	(0, 0, 0, 0, 2, 2)	(0, 0, 0, 0, 0, 0)	$6S$	648	$3G$	$2B$				$6T$
$6D$	9	1	81	(0, 0, 0, 0, 0, 0)	(0, 0, 0, 0, 0, 0)	$6T$	324	$3O$	$2B$				$6K$
		1	81	(0, 0, 0, 0, 1, 1)	(0, 0, 0, 0, 0, 0)	$6U$	324	$3R$	$2B$				$6Y$
		1	81	(0, 0, 0, 0, 2, 2)	(0, 0, 0, 0, 0, 0)	$6V$	324	$3R$	$2B$				$6X$
		2	162	(0, 0, 0, 0, 0, 1)	(0, 0, 0, 0, 0, 0)	$6W$	162	$3P$	$2B$				$6T$
		2	162	(0, 0, 0, 0, 1, 0)	(0, 0, 0, 0, 0, 0)	$6X$	162	$3Q$	$2B$				$6U$
		2	162	(0, 0, 0, 0, 2, 0)	(0, 0, 0, 0, 0, 0)	$6Y$	162	$3S$	$2B$				$6AB$
$6E$	9	1	81	(0, 0, 0, 0, 0, 0)	(0, 0, 0, 0, 0, 0)	$6Z$	324	$3T$	$2B$				$6N$
		1	81	(0, 0, 0, 0, 1, 1)	(0, 0, 0, 0, 0, 0)	$6AA$	324	$3W$	$2B$				$6AA$
		1	81	(0, 0, 0, 0, 2, 2)	(0, 0, 0, 0, 0, 0)	$6AB$	324	$3W$	$2B$				$6Z$
		2	162	(0, 0, 0, 0, 0, 1)	(0, 0, 0, 0, 0, 0)	$6AC$	162	$3U$	$2B$				$6V$
		2	162	(0, 0, 0, 0, 1, 0)	(0, 0, 0, 0, 0, 0)	$6AD$	162	$3V$	$2B$				$6W$
		2	162	(0, 0, 0, 0, 2, 0)	(0, 0, 0, 0, 0, 0)	$6AE$	162	$3X$	$2B$				$6AB$
$8A$	1	1	729	(0, 0, 0, 0, 0, 0)	(0, 0, 0, 0, 0, 0)	$8A$	8	$4B$					$8A$
$9A$	9	1	81	(0, 0, 0, 0, 0, 0)	(0, 0, 0, 0, 0, 0)	$9K$	243		$3D$				$9A$
		1	81	(0, 0, 0, 0, 1, 1)	(0, 0, 0, 0, 0, 0)	$9L$	243		$3D$				$9T$
		1	81	(0, 1, 0, 0, 2, 0)	(0, 0, 0, 0, 0, 0)	$9M$	243		$3D$				$9J$
		3	243	(0, 0, 0, 0, 0, 1)	(0, 0, 0, 0, 0, 0)	$9N$	81		$3J$				$9L$
		3	243	(0, 0, 0, 0, 0, 2)	(0, 0, 0, 0, 0, 0)	$9O$	81		$3G$				$9K$
$9B$	9	1	81	(0, 0, 0, 0, 0, 0)	(0, 0, 0, 0, 0, 0)	$9P$	243		$3D$				$9D$
		1	81	(1, 0, 0, 0, 1, 1)	(0, 0, 0, 0, 0, 0)	$9Q$	243		$3D$				$9H$
		1	81	(2, 0, 0, 0, 2, 0)	(0, 0, 0, 0, 0, 0)	$9R$	243		$3D$				$9G$
		3	243	(0, 0, 0, 0, 1, 0)	(0, 0, 0, 0, 0, 0)	$9S$	81		$3E$				$9N$
		3	243	(0, 0, 0, 0, 2, 0)	(0, 0, 0, 0, 0, 0)	$9T$	81		$3F$				$9M$
$10A$	1	1	729	(0, 0, 0, 0, 0, 0)	(0, 0, 0, 0, 0, 0)	$10A$	20	$5A$					$10A$
$12A$	1	1	729	(0, 0, 0, 0, 0, 0)	(0, 0, 0, 0, 0, 0)	$12G$	36	$6G$	$4A$				$12R$
$12B$	1	1	729	(0, 0, 0, 0, 0, 0)	(0, 0, 0, 0, 0, 0)	$12H$	36	$6I$	$4A$				$12O$
$12C$	9	1	81	(0, 0, 0, 0, 0, 0)	(0, 0, 0, 0, 0, 0)	$12I$	108	$6K$	$4B$				$12A$
		1	81	(0, 0, 0, 0, 0, 1)	(0, 0, 0, 0, 0, 0)	$12J$	108	$6M$	$4B$				$12N$
		1	81	(0, 0, 0, 0, 0, 2)	(0, 0, 0, 0, 0, 0)	$12K$	108	$6L$	$4B$				$12M$
		1	81	(0, 0, 0, 1, 0, 0)	(0, 0, 0, 0, 0, 0)	$12L$	108	$6Q$	$4B$				$12H$
		1	81	(0, 0, 0, 1, 0, 1)	(0, 0, 0, 0, 0, 0)	$12M$	108	$6S$	$4B$				$12K$
		1	81	(0, 0, 0, 1, 0, 2)	(0, 0, 0, 0, 0, 0)	$12N$	108	$6R$	$4B$				$12I$
		1	81	(0, 0, 0, 2, 0, 0)	(0, 0, 0, 0, 0, 0)	$12O$	108	$6N$	$4B$				$12E$
		1	81	(0, 0, 0, 2, 0, 1)	(0, 0, 0, 0, 0, 0)	$12P$	108	$6P$	$4B$				$12F$
		1	81	(0, 0, 0, 2, 0, 2)	(0, 0, 0, 0, 0, 0)	$12Q$	108	$6O$	$4B$				$12L$
$13A$	1	1	729	(0, 0, 0, 0, 0, 0)	(0, 0, 0, 0, 0, 0)	$13A$	13	$13D$		$13D$		$1A$	$13B$
$13B$	1	1	729	(0, 0, 0, 0, 0, 0)	(0, 0, 0, 0, 0, 0)	$13B$	13	$13A$		$13A$		$1A$	$13A$
$13C$	1	1	729	(0, 0, 0, 0, 0, 0)	(0, 0, 0, 0, 0, 0)	$13C$	13	$13B$		$13B$		$1A$	$13B$
$13D$	1	1	729	(0, 0, 0, 0, 0, 0)	(0, 0, 0, 0, 0, 0)	$13D$	13	$13C$		$13C$		$1A$	$13A$
$20A$	1	1	729	(0, 0, 0, 0, 0, 0)	(0, 0, 0, 0, 0, 0)	$20A$	20	$10A$		$4A$			$20A$
$20B$	1	1	729	(0, 0, 0, 0, 0, 0)	(0, 0, 0, 0, 0, 0)	$20B$	20	$10A$		$4A$			$20A$

4. Inertia Factor Groups of $\overline{G} = 3^6:L_4(3)$

We have already seen in Section 3 that the action of $G = L_4(3)$ on $N = 3^6$ has four orbits of lengths 1, 234, 234 and 260. By Brauer's theorem (see Theorem 5.1.5 in [18]), the action of G on $Irr(N)$ will also have four orbits of lengths 1, r , s and t with $1+r+s+t = 729$ such that $[G:H_1] = 1$, $[G:H_2] = r$, $[G:H_3] = s$ and $[G:H_4] = t$, where H_1, H_2, H_3 and H_4 are the inertia factor groups of \overline{G} . When N is an elementary abelian p -group, then it can be regarded as a vector space V over $F = GF(p)$. Any vector space has a dual (the set of all linear functional maps from V into F) denoted by V^* . Although V and V^* are isomorphic as vector spaces (so they have the same dimension), they may not be equivalent as G -modules. It is easy to show that $N^* = Irr(N)$ and hence the action of G on $Irr(N)$ is the same as action of G on N^* . Seretlo [24] developed a programme for the action of G on V^* . It is found in the Brauer ATLAS that $G \cong O_6^+(3)$ has only one irreducible module of dimension 6 over $GF(3)$ and thus N and N^* are equivalent as G -modules. Hence the actions of G on N and N^* are isomorphic, and so the point stabilizers P_i and stabilizers on N^* (inertia factor groups) are the same. We now have that $r = 234$, $s = 234$ and $t = 260$ and that $H_1 = L_4(3)$, $H_2 \cong H_3 = U_4(2)$ and $H_4 = 3^4:2(A_4 \times A_4)$. Using GAP, the number of conjugacy classes of the inertia factor groups are determined and it turns out that

$$|Irr(H_1)| + |Irr(H_2)| + |Irr(H_3)| + |Irr(H_4)| = 29 + 20 + 20 + 42 = 111.$$

This shows that the total contribution of irreducible characters from the four inertia groups is 111 and is equal to the number of classes of \overline{G} as determined in Section 3. The inertia factor groups of \overline{G} are constructed from elements within $G = L_4(3)$ and their generators are as follows:

$H_2 = U_4(2) = \langle \alpha_1, \alpha_2 \rangle$, $\alpha_1 \in 6A$, $\alpha_2 \in 4B$, where

$$\alpha_1 = \begin{pmatrix} 1 & 0 & 1 & 0 & 2 & 0 \\ 1 & 0 & 0 & 0 & 0 & 2 \\ 2 & 0 & 0 & 2 & 2 & 2 \\ 0 & 0 & 0 & 1 & 0 & 0 \\ 1 & 1 & 0 & 1 & 2 & 1 \\ 0 & 0 & 0 & 0 & 0 & 1 \end{pmatrix}, \quad \alpha_2 = \begin{pmatrix} 2 & 2 & 0 & 1 & 0 & 1 \\ 1 & 0 & 1 & 2 & 0 & 1 \\ 2 & 2 & 2 & 2 & 1 & 2 \\ 2 & 2 & 1 & 0 & 0 & 1 \\ 2 & 0 & 2 & 2 & 0 & 2 \\ 0 & 0 & 0 & 0 & 0 & 1 \end{pmatrix}$$

$H_3 = U_4(2) = \langle \beta_1, \beta_2 \rangle$, $\beta_1 \in 6B$, $\beta_2 \in 4C$, where

$$\beta_1 = \begin{pmatrix} 2 & 2 & 2 & 2 & 0 & 0 \\ 1 & 2 & 2 & 1 & 0 & 1 \\ 0 & 1 & 0 & 2 & 1 & 0 \\ 0 & 1 & 1 & 1 & 2 & 1 \\ 0 & 2 & 2 & 0 & 2 & 2 \\ 2 & 0 & 1 & 1 & 2 & 1 \end{pmatrix}, \quad \beta_2 = \begin{pmatrix} 2 & 2 & 1 & 0 & 1 & 1 \\ 1 & 2 & 2 & 0 & 1 & 1 \\ 2 & 0 & 0 & 1 & 0 & 2 \\ 2 & 2 & 2 & 1 & 0 & 1 \\ 1 & 1 & 1 & 0 & 1 & 2 \\ 0 & 2 & 1 & 1 & 2 & 0 \end{pmatrix}$$

$H_4 = 3^4:2(A_4 \times A_4) = \langle \gamma_1, \gamma_2, \gamma_3 \rangle$, $\gamma_1, \gamma_2, \gamma_3 \in 2A$, where

$$\gamma_1 = \begin{pmatrix} 1 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 & 0 & 0 \\ 2 & 1 & 1 & 0 & 2 & 1 \\ 1 & 2 & 2 & 1 & 2 & 2 \\ 2 & 1 & 1 & 2 & 2 & 2 \end{pmatrix}, \quad \gamma_2 = \begin{pmatrix} 1 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 & 0 & 0 \\ 2 & 1 & 1 & 0 & 2 & 1 \\ 1 & 2 & 2 & 1 & 2 & 2 \\ 2 & 1 & 1 & 2 & 2 & 2 \end{pmatrix},$$

$$\gamma_3 = \begin{pmatrix} 1 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 & 0 & 0 \\ 2 & 1 & 1 & 0 & 2 & 1 \\ 1 & 2 & 2 & 1 & 2 & 2 \\ 2 & 1 & 1 & 2 & 2 & 2 \end{pmatrix}$$

5. The fusion of H_2 , H_3 and H_4 into $G = L_4(3)$

The fusion maps of the inertia factor groups H_2 , H_3 and H_4 into $G = L_4(3)$ are obtained, by using the generators of the H_i 's and the GAP command “*FusionConjugacyClasses*(H_i, G)”. The complete fusion maps of H_2 , H_3 and H_4 into $G = L_4(3)$ are shown in Tables 4, 5 and 6 below.

Table 4: The fusion of H_2 into G

$[h]_{H_2} \longrightarrow [g]_{L_4(3)}$	$[h]_{H_2} \longrightarrow [g]_{L_4(3)}$	$[h]_{H_2} \longrightarrow [g]_{L_4(3)}$	$[h]_{H_2} \longrightarrow [g]_{L_4(3)}$
1A 1A	3C 3D	6A 6C	6F 6B
2A 2B	3D 3C	6B 6C	9A 9B
2B 2A	4A 4B	6C 6E	9B 9B
3A 3A	4B 4C	6D 6D	12A 12C
3B 3A	5A 5A	6E 6D	12B 12C

Table 5: The fusion of H_3 into G

$[h]_{H_3} \longrightarrow [g]_{L_4(3)}$	$[h]_{H_3} \longrightarrow [g]_{L_4(3)}$	$[h]_{H_3} \longrightarrow [g]_{L_4(3)}$	$[h]_{H_4} \longrightarrow [g]_{L_4(3)}$
1A 1A	3C 3A	6A 6D	6F 6A
2A 2B	3D 3D	6B 6C	9A 9A
2B 2A	4A 4B	6C 6C	9B 9A
3A 3C	4B 4C	6D 6E	12A 12C
3B 3A	5A 5A	6E 6E	12B 12C

Table 6: The fusion of H_4 into G

$[h]_{H_4} \longrightarrow [g]_{L_4(3)}$	$[h]_{H_4} \longrightarrow [g]_{L_4(3)}$	$[h]_{H_4} \longrightarrow [g]_{L_4(3)}$	$[h]_{H_4} \longrightarrow [g]_{L_4(3)}$
1A 1A	3I 3B	4A 4B	6I 6C
2A 2B	3J 3C	4B 4B	6J 6D
2B 2A	3K 3B	6A 6A	9A 9B
3A 3C	3L 3A	6B 6B	9B 9AC
3B 3A	3M 3B	6C 6C	9C 9B
3C 3D	3N 3D	6D 6C	9D 9A
3D 3A	3O 3B	6E 6E	12A 12C
3E 3B	3P 3A	6F 6D	12B 12C
3F 3A	3Q 3B	6G 6C	12C 12C
3G 3B	3R 3C	6H 6E	12D 12C
3H 3D	3S 3B		

6. The Fischer-Matrices of $\overline{G} = 3^6:L_4(3)$

Having obtained the conjugacy classes of \overline{G} in coset-analysis format and the fusion maps of the inertia factor groups H_2 , H_3 and H_4 into G , the Fischer-Clifford matrices of the group $\overline{G} = 3^6:L_4(3)$ will be now computed. Programme D in [5] is largely used for automatic determination of a possible candidate for each Fischer-Clifford matrix $M(g)$, $g \in G$, of \overline{G} . Then the properties of Fischer-Clifford matrices discussed in detail in [6], [13], [14] and [23] are used to rearrange the rows and columns of this candidate in order to get the unique matrix $M(g)$ for $\overline{G} = 3^6:L_4(3)$. The Programme D only works on split extensions $N:G$, where N is elementary abelian. Note that since $N = 3^6$ is an elementary abelian p -group, then all the relations hold. For example, considering the conjugacy class $2A$ of $G = L_4(3)$, and by making use of Theorem 5.2.4 and property (e) in [18], $M(2A)$ has the following form with corresponding weights attached to the rows and columns.

$$\begin{array}{rccccc}
 & |C_{\overline{G}}(2A)| & |C_{\overline{G}}(6A)| & |C_{\overline{G}}(6B)| & |C_{\overline{G}}(6C)| \\
 & 233280 & 11664 & 7776 & 7776 \\
 |C_{H_1}(2A)| = 2880 & \left(\begin{array}{cccc} 1 & 1 & 1 & 1 \\ 30 & f & g & h \\ 30 & j & k & l \\ 20 & n & o & p \end{array} \right) \\
 |C_{H_2}(2B)| = 96 & & & & \\
 |C_{H_3}(2B)| = 96 & & & & \\
 |C_{H_4}(2B)| = 144 & & & & \\
 m_j & 9 & 180 & 270 & 270
 \end{array}$$

In order to determine the entries f, g, h, j, k, l, n, o and p of the Fischer-Clifford matrix $M(2A)$, the GAP output for programme D for the matrix $M(2A)$ is first generated as,

$$M(2A') = \begin{pmatrix} 1 & 1 & 1 & 1 \\ 20 & 2 & 2 & -7 \\ 30 & 3 & -6 & 3 \\ 30 & -6 & 3 & 3 \end{pmatrix}.$$

Clearly, row 2 of matrix $M(2A')$ becomes row 4 of Fischer-Clifford matrix $M(2A)$. Using the centralizer orders for the classes $6A$, $6B$ and $6C$ of \overline{G} , column 4 of $M(2A')$ becomes column 2 of $M(2A)$ and thus, $f = 3$, $j = 3$ and $n = 7$. Now, columns 3 and 4 of matrix $M(2A)$ are going to be identified using the fact that for any p -singular element g of any finite group G and irreducible character χ of G , then $\chi(g) \equiv \chi(g^p) \pmod{p}$. Noting from Table 3 that $\chi(6B) \pmod{2} \equiv \chi(3A)$, the second power map of the class $6B$ of \overline{G} is applied. If we suppose that column 3 in $M(2A')$ becomes column 3 of $M(2A)$, it turns out that $\chi(6B) \pmod{2} \not\equiv \chi(3A)$ for all irreducible characters of \overline{G} on classes $3A$ and $6B$ of \overline{G} and coming from inertia factor groups H_2 , H_3 and H_4 . Thus column 2 of matrix $M(2A')$ becomes column 3 of $M(2A)$ while column 3 of $M(2A')$ automatically becomes column 4 of $M(2A)$. The final Fischer-Clifford matrix assumes the following structure with $g = 3$, $h = -6$, $k = -6$, $l = 3$, $o = 2$ and $p = 2$.

$$M(2A) = \begin{array}{ccccc} & 233280 & 11664 & 7776 & 7776 \\ \begin{array}{c} 2880 \\ 96 \\ 96 \\ 144 \end{array} & \begin{pmatrix} 1 & 1 & 1 & 1 \\ 30 & 3 & 3 & -6 \\ 30 & 3 & -6 & 3 \\ 20 & -7 & 2 & 2 \end{pmatrix} & & & \\ & 9 & 180 & 270 & 270 \end{array}$$

With a quite a number of the Fischer-Clifford matrices of \overline{G} containing entries which involve the complex number $\alpha = -\frac{1}{2} + \frac{\sqrt{-3}}{2}$ such that $\alpha^3 = 1$, the Fischer-Clifford matrix corresponding to the conjugacy class $6D$ of $G = L_4(3)$ is briefly discussed. Using theorem 5.2.4 and property (e) in [18], $M(6D)$ has the following form with corresponding weights attached to the rows and columns:

$$\begin{array}{l} |C_{\overline{G}}(6T)| \quad |C_{\overline{G}}(6U)| \quad |C_{\overline{G}}(6V)| \quad |C_{\overline{G}}(6W)| \quad |C_{\overline{G}}(6X)| \quad |C_{\overline{G}}(6Y)| \\ \begin{array}{cccccc} 324 & 324 & 324 & 162 & 162 & 162 \end{array} \\ \begin{array}{l} |C_{H_1}(6D)| = 36 \\ |C_{H_2}(6D)| = 36 \\ |C_{H_3}(6E)| = 36 \\ |C_{H_4}(6A)| = 18 \\ |C_{H_4}(6F)| = 18 \\ |C_{H_4}(6J)| = 18 \end{array} \left(\begin{array}{cccccc} 1 & 1 & 1 & 1 & 1 & 1 \\ 1 & a & f & k & p & u \\ 1 & b & g & l & q & v \\ 2 & c & h & m & r & w \\ 2 & d & i & n & s & x \\ 2 & e & j & o & t & y \end{array} \right) \\ \begin{array}{cccccc} m_j & 81 & 81 & 81 & 162 & 162 & 162 \end{array} \end{array}$$

To determine the remaining entries of the Fischer-Clifford matrix $M(6D)$, the GAP output for programme D for $M(6D)$ is computed,

$$M(6D') = \begin{pmatrix} 1 & 1 & 1 & 1 & 1 & 1 \\ 2 & 2 & 2 & -1 & -1 & -1 \\ 2 & 2\bar{\alpha} & 2\alpha & -\alpha & -1 & -\bar{\alpha} \\ 1 & \bar{\alpha} & \alpha & \alpha & 1 & -\bar{\alpha} \\ 2 & 2\alpha & 2\bar{\alpha} & -\bar{\alpha} & -1 & -\alpha \\ 1 & \alpha & \bar{\alpha} & \bar{\alpha} & 1 & \alpha \end{pmatrix}$$

where $\bar{\alpha} = -\frac{1}{2} - \frac{\sqrt{-3}}{2}$ and $\alpha = -\frac{1}{2} + \frac{\sqrt{-3}}{2}$. The rows of Fischer-Clifford matrix $M(6D')$ are then rearranged to match the structure above. Thus we have

$$M(6D') = \begin{pmatrix} 1 & 1 & 1 & 1 & 1 & 1 \\ 1 & \bar{\alpha} & \alpha & \alpha & 1 & -\bar{\alpha} \\ 1 & \alpha & \bar{\alpha} & \bar{\alpha} & 1 & \alpha \\ 2 & 2 & 2 & -1 & -1 & -1 \\ 2 & 2\bar{\alpha} & 2\alpha & -\alpha & -1 & -\bar{\alpha} \\ 2 & 2\alpha & 2\bar{\alpha} & -\bar{\alpha} & -1 & -\alpha \end{pmatrix}$$

Taking into consideration properties of Fischer-Clifford matrices discussed in [6], [13], [14] and [23], the values of the irreducible characters of \overline{G} on the classes $3R$, $3P$, $3Q$ and $3S$ that have already been obtained by Fischer-Clifford matrix $M(3C)$, and using the fact that $(6U)^2 = 3R$, $(6V)^2 = 3R$, $(6W)^2 = 3P$, $(6X)^2 = 3Q$ and $(6Y)^2 = 3S$ (see Table 3), the final Fischer-Clifford matrix $M(6D)$ assumes the form below.

$$M(6D) = \begin{array}{cccccc} & 324 & 324 & 324 & 162 & 162 & 162 \\ \begin{array}{l} 36 \\ 36 \\ 36 \\ 18 \\ 18 \\ 18 \end{array} & \begin{pmatrix} 1 & 1 & 1 & 1 & 1 & 1 \\ 1 & \alpha & \bar{\alpha} & \alpha & \bar{\alpha} & 1 \\ 1 & \bar{\alpha} & \alpha & \bar{\alpha} & \alpha & 1 \\ 2 & 2 & 2 & -1 & -1 & -1 \\ 2 & 2\bar{\alpha} & 2\alpha & -\bar{\alpha} & -\alpha & -1 \\ 2 & 2\alpha & 2\bar{\alpha} & -\alpha & -\bar{\alpha} & -1 \end{pmatrix} \\ & 81 & 81 & 81 & 162 & 162 & 162 \end{array}$$

For each class representative $g \in L_4(3)$, a Fischer-Clifford matrix $M(g)$ is constructed and listed in Table 7.

TABLE 7. The Fischer-Clifford Matrices of $\bar{G} = 3^6:L_4(3)$

$M(g)$	$M(g)$
$M(1A) = \begin{pmatrix} 1 & 1 & 1 & 1 \\ 234 & -9 & 18 & -9 \\ 234 & 18 & -9 & -9 \\ 260 & -10 & -10 & 17 \end{pmatrix}$	$M(2A) = \begin{pmatrix} 1 & 1 & 1 & 1 \\ 30 & 3 & 3 & -6 \\ 30 & 3 & -6 & 3 \\ 20 & -7 & 2 & 2 \end{pmatrix}$
$M(2B) = \begin{pmatrix} 1 & 1 & 1 & 1 \\ 2 & 2 & -1 & -1 \\ 2 & -1 & 2 & -1 \\ 4 & -2 & -2 & 1 \end{pmatrix}$	$M(3A) = \begin{pmatrix} 1 & 1 & 1 & 1 & 1 & 1 & 1 & 1 & 1 & 1 \\ 9 & 9 & 9 & 9\bar{\alpha} & 9\bar{\alpha} & 9\bar{\alpha} & 9\alpha & 9\alpha & 9\alpha & 0 \\ 9 & 9 & 9 & 9\alpha & 9\alpha & 9\alpha & 9\bar{\alpha} & 9\bar{\alpha} & 9\bar{\alpha} & 0 \\ 9 & 9\bar{\alpha} & 9\alpha & 9 & 9\bar{\alpha} & 9\alpha & 9 & 9\bar{\alpha} & 9\alpha & 0 \\ 9 & 9\alpha & 9\bar{\alpha} & 9 & 9\alpha & 9\bar{\alpha} & 9 & 9\alpha & 9\bar{\alpha} & 0 \\ 8 & 8 & 8 & 8 & 8 & 8 & 8 & 8 & 8 & -1 \\ 9 & 9\bar{\alpha} & 9\alpha & 9\bar{\alpha} & 9\alpha & 9 & 9\alpha & 9 & 9\bar{\alpha} & 0 \\ 9 & 9\alpha & 9\bar{\alpha} & 9\alpha & 9\bar{\alpha} & 9 & 9\bar{\alpha} & 9 & 9\alpha & 0 \\ 9 & 9\bar{\alpha} & 9\alpha & 9\alpha & 9 & 9\bar{\alpha} & 9\bar{\alpha} & 9\alpha & 9 & 0 \\ 9 & 9\alpha & 9\bar{\alpha} & 9\bar{\alpha} & 9 & 9\alpha & 9\alpha & 9\bar{\alpha} & 9 & 0 \end{pmatrix}$
$M(3B) = \begin{pmatrix} 1 & 1 & 1 & 1 & 1 & 1 & 1 & 1 & 1 \\ 1 & 1 & 1 & \bar{\alpha} & \bar{\alpha} & \bar{\alpha} & \alpha & \alpha & \alpha \\ 1 & 1 & 1 & \alpha & \alpha & \alpha & \bar{\alpha} & \bar{\alpha} & \bar{\alpha} \\ 1 & \bar{\alpha} & \alpha & 1 & \bar{\alpha} & \alpha & 1 & \bar{\alpha} & \alpha \\ 1 & \bar{\alpha} & \alpha & \bar{\alpha} & \alpha & 1 & \alpha & 1 & \bar{\alpha} \\ 1 & \bar{\alpha} & \alpha & \alpha & 1 & \bar{\alpha} & \bar{\alpha} & \alpha & 1 \\ 1 & \alpha & \bar{\alpha} & 1 & \alpha & \bar{\alpha} & 1 & \alpha & \bar{\alpha} \\ 1 & \alpha & \bar{\alpha} & \bar{\alpha} & 1 & \alpha & \alpha & \bar{\alpha} & 1 \\ 1 & \alpha & \bar{\alpha} & \alpha & \bar{\alpha} & 1 & \bar{\alpha} & 1 & \alpha \end{pmatrix}$	$M(3C) = \begin{pmatrix} 1 & 1 & 1 & 1 & 1 & 1 \\ 18 & -9 & -9 & 9 & 0 & 0 \\ 36 & 9 & 9 & 0 & -9 & 0 \\ 2 & 2 & 2 & 2 & 2 & -1 \\ 12 & 6\alpha - 3\bar{\alpha} & -3\alpha + 6\bar{\alpha} & -6 & 3 & 0 \\ 12 & -3\alpha + 6\bar{\alpha} & 6\alpha - 3\bar{\alpha} & -6 & 3 & 0 \end{pmatrix}$
$M(3D) = \begin{pmatrix} 1 & 1 & 1 & 1 & 1 & 1 \\ 18 & -9 & -9 & 9 & 0 & 0 \\ 36 & 9 & 9 & 0 & -9 & 0 \\ 2 & 2 & 2 & 2 & 2 & -1 \\ 12 & 6\alpha - 3\bar{\alpha} & -3\alpha + 6\bar{\alpha} & -6 & 3 & 0 \\ 12 & -3\alpha + 6\bar{\alpha} & 6\alpha - 3\bar{\alpha} & -6 & 3 & 0 \end{pmatrix}$	$M(4A) = (1)$
$M(4B) = \begin{pmatrix} 1 & 1 & 1 & 1 & 1 \\ 2 & 2 & -1 & -1 & -1 \\ 2 & -1 & -1 & 2 & -1 \\ 2 & -1 & 2 & -1 & -1 \\ 2 & -1 & -1 & -1 & 2 \end{pmatrix}$	$M(4C) = \begin{pmatrix} 1 & 1 & 1 \\ 4 & -2 & 1 \\ 4 & 1 & -2 \end{pmatrix}$
$M(5A) = \begin{pmatrix} 1 & 1 & 1 \\ 4 & 1 & -2 \\ 4 & -2 & 1 \end{pmatrix}$	$M(6A) = \begin{pmatrix} 1 & 1 & 1 \\ 6 & -3 & 0 \\ 2 & 2 & -1 \end{pmatrix}$

$$\bar{\alpha} = -\frac{1}{2} - \frac{\sqrt{-3}}{2}, \alpha = -\frac{1}{2} + \frac{\sqrt{-3}}{2}$$

TABLE 7. (continued)

$M(g)$	$M(g)$
$M(6B) = \begin{pmatrix} 1 & 1 & 1 \\ 6 & -3 & 0 \\ 2 & 2 & -1 \end{pmatrix}$	$M(6C) = \begin{pmatrix} 1 & 1 & 1 & 1 & 1 & 1 & 1 & 1 & 1 \\ 1 & 1 & 1 & \bar{\alpha} & \bar{\alpha} & \bar{\alpha} & \alpha & \alpha & \alpha \\ 1 & 1 & 1 & \alpha & \alpha & \alpha & \bar{\alpha} & \bar{\alpha} & \bar{\alpha} \\ 1 & \bar{\alpha} & \alpha & \alpha & 1 & \bar{\alpha} & \bar{\alpha} & \alpha & 1 \\ 1 & \alpha & \bar{\alpha} & \bar{\alpha} & 1 & \alpha & \alpha & \bar{\alpha} & 1 \\ 1 & \alpha & \bar{\alpha} & 1 & \alpha & \bar{\alpha} & 1 & \alpha & \bar{\alpha} \\ 1 & \bar{\alpha} & \alpha & 1 & \bar{\alpha} & \alpha & 1 & \bar{\alpha} & \alpha \\ 1 & \alpha & \bar{\alpha} & \alpha & \bar{\alpha} & 1 & \bar{\alpha} & 1 & \alpha \\ 1 & \bar{\alpha} & \alpha & \bar{\alpha} & \alpha & 1 & \alpha & 1 & \bar{\alpha} \end{pmatrix}$
$M(6D) = \begin{pmatrix} 1 & 1 & 1 & 1 & 1 & 1 \\ 1 & \alpha & \bar{\alpha} & \alpha & \bar{\alpha} & 1 \\ 1 & \bar{\alpha} & \alpha & \bar{\alpha} & \alpha & 1 \\ 2 & 2 & 2 & -1 & -1 & -1 \\ 2 & 2\bar{\alpha} & 2\alpha & -\bar{\alpha} & -\alpha & -1 \\ 2 & 2\alpha & 2\bar{\alpha} & -\alpha & -\bar{\alpha} & -1 \end{pmatrix}$	$M(6E) = \begin{pmatrix} 1 & 1 & 1 & 1 & 1 & 1 \\ 2 & 2 & 2 & -1 & -1 & -1 \\ 1 & \alpha & \bar{\alpha} & \alpha & \bar{\alpha} & 1 \\ 1 & \bar{\alpha} & \alpha & \bar{\alpha} & \alpha & 1 \\ 2 & 2\bar{\alpha} & 2\alpha & -\bar{\alpha} & -\alpha & -1 \\ 2 & 2\alpha & 2\bar{\alpha} & -\alpha & -\bar{\alpha} & -1 \end{pmatrix}$
$M(8A) = (1)$	$M(9A) = \begin{pmatrix} 1 & 1 & 1 & 1 & 1 \\ 3 & 3\bar{\alpha} & 3\alpha & 0 & 0 \\ 3 & 3\alpha & 3\bar{\alpha} & 0 & 0 \\ 1 & 1 & 1 & \bar{\alpha} & \alpha \\ 1 & 1 & 1 & \alpha & \bar{\alpha} \end{pmatrix}$
$M(9B) = \begin{pmatrix} 1 & 1 & 1 & 1 & 1 \\ 3 & 3\bar{\alpha} & 3\alpha & 0 & 0 \\ 3 & 3\alpha & 3\bar{\alpha} & 0 & 0 \\ 1 & 1 & 1 & \bar{\alpha} & \alpha \\ 1 & 1 & 1 & \alpha & \bar{\alpha} \end{pmatrix}$	$M(10A) = (1)$
$M(12A) = (1)$	$M(12B) = (1)$
$M(12C) = \begin{pmatrix} 1 & 1 & 1 & 1 & 1 & 1 & 1 & 1 \\ 1 & 1 & 1 & \bar{\alpha} & \bar{\alpha} & \bar{\alpha} & \alpha & \alpha \\ 1 & 1 & 1 & \alpha & \alpha & \alpha & \bar{\alpha} & \bar{\alpha} \\ 1 & \alpha & \bar{\alpha} & \bar{\alpha} & 1 & \alpha & \alpha & \bar{\alpha} \\ 1 & \bar{\alpha} & \alpha & \alpha & 1 & \bar{\alpha} & \bar{\alpha} & \alpha \\ 1 & \bar{\alpha} & \alpha & 1 & \bar{\alpha} & \alpha & 1 & \bar{\alpha} \\ 1 & \alpha & \bar{\alpha} & 1 & \alpha & \bar{\alpha} & 1 & \alpha \\ 1 & \bar{\alpha} & \alpha & \bar{\alpha} & \alpha & 1 & \alpha & 1 \\ 1 & \alpha & \bar{\alpha} & \alpha & \bar{\alpha} & 1 & \bar{\alpha} & 1 \end{pmatrix}$	$M(13A) = (1)$
$M(13B) = (1)$	$M(13C) = (1)$
$M(13D) = (1)$	$M(20A) = (1)$
$M(20B) = (1)$	

$$\bar{\alpha} = -\frac{1}{2} - \frac{\sqrt{-3}}{2}, \alpha = -\frac{1}{2} + \frac{\sqrt{-3}}{2}$$

7. The Character Table of $\bar{G} = 3^6:L_4(3)$

Having obtained the conjugacy classes of $\bar{G} = 3^6:L_4(3)$, the ordinary character tables of all the inertia factor groups available in GAP [11], the fusions of conjugacy classes of the inertia factors into classes of $L_4(3)$ and the Fischer-Clifford matrices of $\bar{G} = 3^6:L_4(3)$, the full character table of \bar{G} can now be constructed by following the theoretical outline discussed in Section 2. This character table is a 111×111 complex-valued matrix partitioned row-wise into four blocks $\Delta_1, \Delta_2, \Delta_3$ and Δ_4 , where each block corresponds to an inertia group $\bar{H}_i = 3^6:H_i$. In fact, $\Delta_1 = \{\chi_i | 1 \leq i \leq 29\}$, $\Delta_2 = \{\chi_i | 30 \leq i \leq 49\}$, $\Delta_3 = \{\chi_i | 50 \leq i \leq 69\}$ and $\Delta_4 = \{\chi_i | 70 \leq i \leq 111\}$, where $\text{Irr}(3^6:L_4(3)) = \cup_{i=1}^4 \Delta_i$. The consistency and accuracy of this table has been tested using Programme C in [22] and the complete character table of the group \bar{G} can be accessed via the link below.

<https://drive.google.com/file/d/16oDwruWA0qWOFJNy8JmgOoESw7UoumVR/view?usp=sharing>

8. The Fusion of $3^6:L_4(3)$ of into $O_8^+(3)$

Since $\bar{G} = 3^6:L_4(3)$ is a maximal subgroup of $O_8^+(3)$ of index 1120, then the action of $O_8^+(3)$ on the cosets of \bar{G} gives rise to a permutation character $\chi(O_8^+(3)|\bar{G})$ of degree 1120. From the ATLAS of finite groups [8], $\chi(O_8^+(3)|\bar{G}) = 1a + 300a + 819a$, where $1a$, $300a$ and $819a$ are irreducible characters of $O_8^+(3)$ of degrees 1, 300 and 819 respectively. Using the information provided by the conjugacy classes of $3^6:L_4(3)$ and $O_8^+(3)$, the power maps and the permutation character of $O_8^+(3)$ of degree 1120, a partial fusion of \bar{G} into $O_8^+(3)$ is obtained. To complete the fusion maps, irreducible characters of $O_8^+(3)$ of small degrees are restricted to $3^6:L_4(3)$. To determine the restrictions of irreducible characters of $O_8^+(3)$ to $3^6:L_4(3)$, the technique of set intersections for characters which has been discussed in detail in [1], [17] and [18] is used.

Let ρ be the character of $L_4(3)$ afforded by the regular representation of $L_4(3)$. Then $\rho = \sum_{i=1}^{29} \varepsilon_i \phi_i$, where $\phi_i \in \text{Irr}(L_4(3))$ and $\varepsilon_i = \deg(\phi_i)$. Thus ρ can be regarded as a character of $3^6:L_4(3)$ which contains 3^6 in its kernel such that,

$$\rho(g) = \begin{cases} |L_4(3)| & \text{if } g \in 3^6 \\ 0 & \text{otherwise} \end{cases}$$

If ψ is a character of $O_8^+(3)$, then

$$\begin{aligned}
\langle \rho, \psi \rangle_{\overline{G}} &= \frac{1}{|3^6:L_4(3)|} \{ \rho(1A)\psi(1A) + 234\rho(3A)\psi(3A) + 234\rho(3B)\psi(3B) + 260\rho(3C)\psi(3C) \} \\
&= \frac{1}{|3^6:L_4(3)|} \{ |L_4(3)|\psi(1A) + 234|L_4(3)|\psi(3A) + 234|L_4(3)|\psi(3B) + 260|L_4(3)|\psi(3C) \} \\
&= \frac{1}{729} \{ \psi(1A) + 234\psi(3A) + 234\psi(3B) + 260\psi(3C) \} \\
&= \langle \psi_{3^6}, \tau_1 \rangle,
\end{aligned}$$

where τ_1 is the identity character of 3^6 and ψ_{3^6} is the restriction of ψ to 3^6 . Also for ψ it is obtained that;

$$\psi_{3^6} = a_1\theta_1 + a_2\theta_2 + a_3\theta_3 + a_4\theta_4$$

where $a_1, a_2, a_3, a_4 \in \mathbf{N} \cup \{0\}$ and $\theta_i, i \in \{1, 2, 3, 4\}$ are sums of the irreducible characters of 3^6 which are in one orbit under action of $L_4(3)$ on $Irr(3^6)$. Letting $\tau_j \in Irr(3^6)$, where $j \in \{1, 2, \dots, 729\}$, then,

$$\begin{aligned}
\theta_1 &= \tau_1, \deg(\theta_1) = 1 \\
\theta_2 &= \sum_{j=2}^{235} \tau_j, \deg(\theta_2) = 234 \\
\theta_3 &= \sum_{j=236}^{469} \tau_j, \deg(\theta_3) = 234 \\
\theta_4 &= \sum_{j=470}^{729} \tau_j, \deg(\theta_4) = 260
\end{aligned}$$

Hence

$$\psi_{3^6} = a_1\tau_1 + a_2 \sum_{j=2}^{235} \tau_j + a_3 \sum_{j=236}^{469} \tau_j + a_4 \sum_{j=470}^{729} \tau_j$$

and,

$$\begin{aligned}
\langle \psi_{3^6}, \psi_{3^6} \rangle &= a_1^2 + 234a_2^2 + 234a_3^2 + 260a_4^2 \\
&= \frac{1}{729} \{ \psi(1A)\psi(1A) + 234\psi(3A)\psi(3A) + 234\psi(3B)\psi(3B) + 260\psi(3C)\psi(3C) \},
\end{aligned}$$

where

$$a_1 = \langle \psi_{3^6}, \tau_1 \rangle = \langle \rho, \psi \rangle_{\overline{G}}$$

Now, the above results are applied to some irreducible characters of $O_8^+(3)$ of small degrees, which in this case are $\psi_1 = 260a$, $\psi_2 = 260b$, $\psi_3 = 260c$, $\psi_4 = 819b$, $\psi_5 = 819c$, $\psi_6 = 2275a$, $\psi_7 = 2275b$, $\psi_8 = 2275c$, and $\psi_9 = 2275d$ of degrees 1, 260, 260, 260, 819, 819, 2275, 2275, 2275, and 2275, respectively (see ATLAS for character table of $O_8^+(3)$). From the partial fusion that has already been determined, the class $3C$ of \overline{G} must fuse into the class $3A$ of $O_8^+(3)$. Also, the classes $3A$ and $3B$ of \overline{G} must each fuse into one of the following classes, $3B$, $3C$, $3D$, $3E$, $3F$ and $3G$ of $O_8^+(3)$ such that the condition $a_1 = \langle \psi_{3^6}, \tau_1 \rangle \in \mathbf{N} \cup \{0\}$ is satisfied. The values of $\psi_1 = 260a$ on the classes $3A$, $3B$ and $3C$ of $O_8^+(3)$ violate this condition. The only combination that satisfies this condition is $3A$, $3B$ and $3E$ and thus for ψ_1 the following result is obtained,

$$a_1 = \langle \rho, \psi \rangle_{\overline{G}} = \frac{1}{729} \{260 + 234(17) + 234(44) + 260(17)\} = 26$$

Since the degree of ψ_1 is 260 then,

$$a_1 + 234a_2 + 234a_3 + 260a_4 = 260$$

and so we must have that $a_1 = 26$, $a_2 = 1$, $a_3 = 0$ and $a_4 = 0$ or $a_1 = 26$, $a_2 = 0$, $a_3 = 1$ and $a_4 = 0$. It turns out that $(\psi_1)_{3^6:L_4(3)}$ is a sum of two irreducible characters of $3^6:L_4(3)$ of degrees 26 and 234, respectively. Based on the partial fusion of $3^6:L_4(3)$ into $O_8^+(3)$ that has already been done, the result below is obtained,

$$(\psi_1)_{3^6:L_4(3)} = \chi_3 + \chi_{50}$$

Similarly, for $\psi_2 = 260b$, a_1 is obtained as follows,

$$a_1 = \langle \rho, \psi \rangle_{\overline{G}} = \frac{1}{729} \{260 + 234(-10) + 234(-10) + 260(17)\} = 0$$

Since the degree of ψ_2 is 260 then

$$a_1 + 234a_2 + 234a_3 + 260a_4 = 260$$

so that $a_1 = 0$, $a_2 = 0$, $a_3 = 0$ and $a_4 = 1$ and thus $(\psi_2)_{3^6:L_4(3)}$ is an irreducible character of $3^6:L_4(3)$ of degree 260. Based on the partial fusion of $3^6:L_4(3)$ into $O_8^+(3)$ that has already been done, it turns out that

$$(\psi_2)_{3^6:L_4(3)} = \chi_{71}.$$

Applying the same procedure to ψ_3 , ψ_4 , ψ_5 , ψ_6 , ψ_7 , ψ_8 and ψ_9 , we obtained the restricted characters below.

$$\begin{aligned} (\psi_3)_{3^6:L_4(3)} &= \chi_{76} \\ (\psi_4)_{3^6:L_4(3)} &= \chi_4 + \chi_{79} \\ (\psi_5)_{3^6:L_4(3)} &= \chi_4 + \chi_{80} \\ (\psi_6)_{3^6:L_4(3)} &= \chi_6 + \chi_{32} + \chi_{90} \\ (\psi_7)_{3^6:L_4(3)} &= \chi_6 + \chi_{31} + \chi_{91} \\ (\psi_8)_{3^6:L_4(3)} &= \chi_7 + \chi_{51} + \chi_{87} \\ (\psi_9)_{3^6:L_4(3)} &= \chi_7 + \chi_{52} + \chi_{86} \end{aligned}$$

By making use of the partial fusion which has already been determined, the values of ψ_1, \dots, ψ_9 on the classes of $O_8^+(3)$ and the values of $(\psi_1)_{3^6:L_4(3)}, \dots, (\psi_9)_{3^6:L_4(3)}$ on the classes of $3^6:L_4(3)$, the fusion of $3^6:L_4(3)$ into $O_8^+(3)$ is completed and is found in Table 3.

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